

(-)

/

-

.(conservative variables)

%

%

)

μm

.(

(0.6kg/s)

()

.()

(. - .)

(-)

(. - .)

(. - .)

(. - .)

STUDY THE EFFECT OF LOADING RATIO ON SECONDARY LOSS AND HEAT EXCHANGE IN TWO PHASE FLOW (GAS – SOLID)IN VARIABLE CROSS SECTIONAL PIPE.

Zaid M. Dulaimy
Technical College Najaf
F.T.E

ABSTRACT

The flow of gas with a suspension of solid particles through venture has been analyzed theoretically. Computations of the flow parameters have been performed based on one-dimensional analysis and numerical solution of the governing equations by (physico – computational model). An investigation of pressure variation, velocity and temperature distributions across venture for air only and for gas-solid suspensions were carried out for specific air mass flow rate and solid loading ratio. Venture was used to produce a non-equilibrium two-phase flow region and the effects of solid particles on the pressure distribution have been considered. The effects of different loading ratios of (0.5 ,2 ,4) on pressure drop, gas and solid velocity and temperature distributions for the same particle diameter of 100 microns and for constant gas mass flow rate of (0.6 kg/s) are determined.

Two sets of diameters of the solid particles (50, 200 microns) were used at the loading rate of (4). The most important conclusions obtained are that the increase loss of pressure due to the increased rate of loading, especially in the convergence of the Venturi, which ranges between (0.7 - 1.3) greater than the decrease of pressure in the divergence of the tube, or a decrease in the size of the particles of solid material, it range (0.7 - 1.1) for a range of particle size of (50-200 microns), where the increase in pressure loss resulting from the decrease of solid particles is greater than the loss of pressure due to increase the loading rate. Increasing the proportion of loading note a decline in the speed of the solid and the change of speed ranges between 1.1 and 2.1 and the change is evident in the spacing of the Venturi. And we get the same effect in the case of decrease in the volume of solid. And shows the amount of change in the speed of solid particles to the change of the size of the particles at a range of (1.6 to 2.7) and appears in the throat and the divergence clearly compared to the convergence of the Venturi.

m^2	A
$\text{kJ / kg} \cdot \text{K}$	C_D
μm	C
	dp

kJ / s		$\dot{\Delta E}_p$
		f
N	()	F_{wx}
m/s^2		g
kJ / kg		h
W / m . K		K
		k
m		L
kg		m
kg / s		\dot{m}
N		ΔM_p
N		$\dot{\Delta M}_A$
		Nu
kN / m^2		p
kN / m^2		Δp_z
		Pr
kJ / s		\dot{Q}
		Re
J / kg . K		R
K		T
s		t
m / s		u
m / s		Ur
m / s		v
kJ / s		w
		z

	Δ
kg / m ³	ρ
N . s / m ²	μ
s	τ
	β
	ϵ
deg	θ
	0
	1
	2
	g
	in
	out
	p
()	th
	z

(Sharma,1978)

(pneumatic conveying system)

(Lee, 1982)

(single – phase flow)

(multiphase flow)

(Govier, 1972)

G.w. Govier & K. Aziz

(Govier, 1972)

shape factor

()

(Efstathios, 1984)

(Navier-Stokes)

(Lee, 1986)

(1.2) (1.06)

(.003 m) (0.0001 m)

(11.8 m/s) (0.53 m/s)

(Stokes drag law)

$$CD = \frac{P_p}{\left(\frac{1}{2} \rho_f U_r^2\right) \frac{1}{4} \pi d_p^2}$$

(Bollio et. Al ,1996)

(Lun & Lin, 2004)

(Osama, 2005)

)

(39.5kg/s)

(μm)

:

(physico – computational model) [1]

(conservative variables)

(source term)

:

.(control volume)

() (&)
 :(Zucker et. al, 2002)

$$\rho_2 u_2 A_2 = \rho_1 u_1 A_1 + \Delta \dot{m}_p \quad (1)$$

(Zucker et. al, 2002)

$$\rho_2 u_2^2 A_2 - \rho_1 u_1^2 A_1 = +p_1 A_1 - p_2 A_2 + \left(\frac{p_1 + p_2}{2} \right) (A_2 - A_1) + \Delta M_p - f_{wx} \quad (2)$$

Zucker et. al,

:(2002)

$$\dot{Q} - \dot{W} = \rho_2 u_2 A_2 \left(h_2 + \frac{u_2^2}{2} \right) - \rho_1 u_1 A_1 \left(h_1 + \frac{u_1^2}{2} \right) - \Delta \dot{E}_p \quad (3)$$

$$\rho u A, p A, \rho u^2 A \quad \& \quad \rho u A \left[h + \left(\frac{u^2}{2} \right) \right]$$

$$\dot{m} = \rho u A \tag{4}$$

$$\text{Impulse function (F)} = pA + \rho u^2 A = pA + \dot{m} u \tag{5}$$

$$h_o = \left(h + \frac{u^2}{2} \right) \rho u A \tag{6}$$

$$h = C_g T_g = \left(\frac{k}{k-1} \right) \left(\frac{p}{\rho} \right) \tag{7}$$

$$: \tag{6} \tag{7}$$

$$h_o = \dot{m} \left[\left(\frac{k}{k-1} \right) \frac{p}{\rho} + \frac{u^2}{2} \right] \tag{8}$$

$$:(u) \tag{5}$$

$$u = \left(\frac{k}{k+1} \right) \left(\frac{F}{\dot{m}} \right) \left[1 \pm \sqrt{1 + 2 \left(\frac{1-k^2}{k^2} \right) \left(\frac{\dot{m} h_o}{F^2} \right)} \right] \tag{9}$$

$$() \tag{}$$

:

$$\Delta \dot{M}_p = \dot{m}_p (v_1 - v_2) = z \dot{m}_g (v_1 - v_2) \tag{10}$$

(Govier, 1972)

$$m_p \frac{dv}{dt} = C D_p (u - v) |u - v| \frac{A_p}{2} + m_p g \tag{11}$$

()

[1]

$$Re = \frac{\rho |u - v| d}{\mu} \tag{12}$$

: (Lee. S.L., 1987) ()

$$C D = \left(\frac{24}{Re} \right) (1 + .15 Re^{0.687}) \tag{13}$$

: () ()

$$\frac{dv}{dt} = \frac{18 \mu f}{\rho_p d^2} (u - v) + g \tag{14}$$

:

$$f = C D \left(\frac{Re}{24} \right) \tag{15}$$

:

$$\tau_d = \frac{\rho_p d^2}{18 \mu f} \tag{16}$$

:

$$\frac{d v}{d t} = \frac{1}{\tau_d} (u - v) + g \tag{17}$$

(&)

:

$$v_2 = u_2 + g \tau_d - \left[((u_1 - v_1) + g \tau_d) \exp\left(-\left(\frac{\Delta x}{v_1 \tau_d}\right)\right) - \frac{\Delta u}{\Delta x} v_1 \tau_d \left(1 - \exp\left(-\frac{\Delta x}{v_1 \tau_d}\right)\right) \right] \tag{18}$$

:

Govier,] (&)

:[1972

$$\Delta E_p = z \dot{m} g \left(\frac{v_1^2 - v_2^2}{2} \right) + z \dot{m} g C_p (t_{p1} - t_{p2}) \tag{19}$$

:[Govier, 1972]

$$m_p \frac{d (c p t_p)}{d t} = h_c (4 \pi d^2) (t_g - t_p) \tag{20}$$

Govier,]

:[1972

$$\frac{d t_p}{d t} = \left(6 N u \frac{k}{d^2 \rho_p c p} \right) (t_g - t_p) \tag{21}$$

[Bolio et. al, 1996]

$$\tau_g = 3 \left(\frac{p r}{N u} \right) \left(\frac{c p}{c g} \right) f \tau_d \tag{22}$$

: [Bolio et. al, 1996]

$$\frac{dt_p}{dt} = \frac{t_g - t_p}{\tau_g} \tag{23}$$

:

$$t_{p2} = t_{g2} + (t_{g1} - t_{p1}) \exp\left(-\frac{\Delta x}{v_1 t_g}\right) - \frac{\Delta t_g}{\Delta x} v_1 t_g \left[1 - \exp\left(-\frac{\Delta x}{v_1 t_g}\right)\right] \tag{24}$$

[Balasim, 2005]

$$Nu = 2 + 0.459 Re^{0.55} pr^{0.33} \tag{25}$$

(Re_p < 1000)

(pr < 0.7)

:

[Holman, 2007]

R= 0.287 (kJ/kg.K) .

g = 9.81 (m/s²) .

μ = 18.6 (N.s/m²)

k = 0.024(kJ/m.s.K)

:

Cp = 0.903 (kJ/kg.s) .

dp= 100 μm

$$\rho_p = 2450 \text{ (kg/m}^3\text{)}$$

:

:

()

(100 μ m)

()

(0.6kg/s)

:

()

()

)

(%)

(0.6 kg / s)

()

.(%

(stokes number)

()

()

[Balasim, 2005]

()

[Balasim, 2005]

.()

()

()

(.)

()

:

 (p / p_0) (z)

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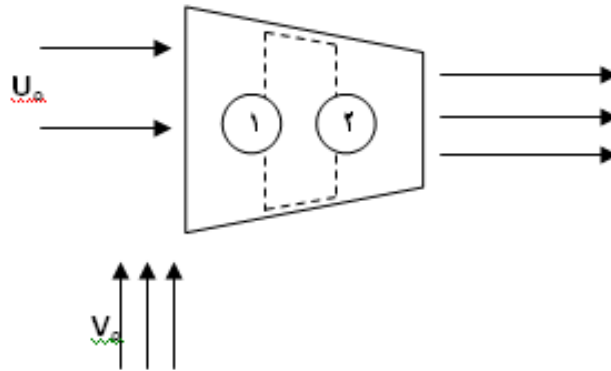
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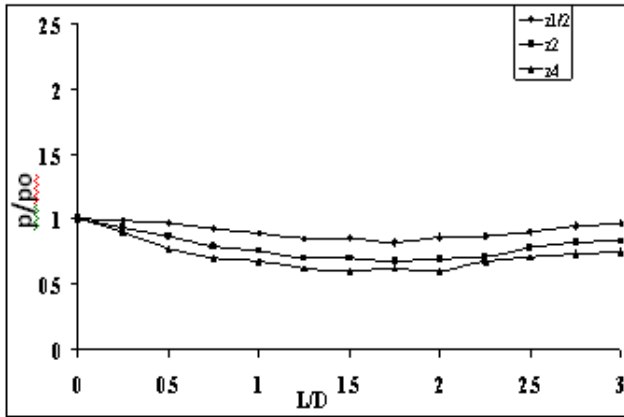
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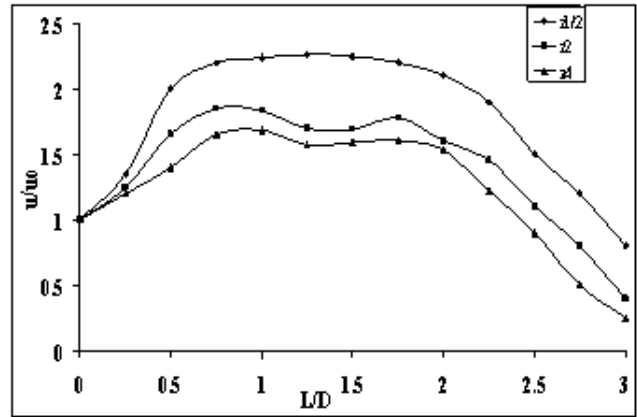
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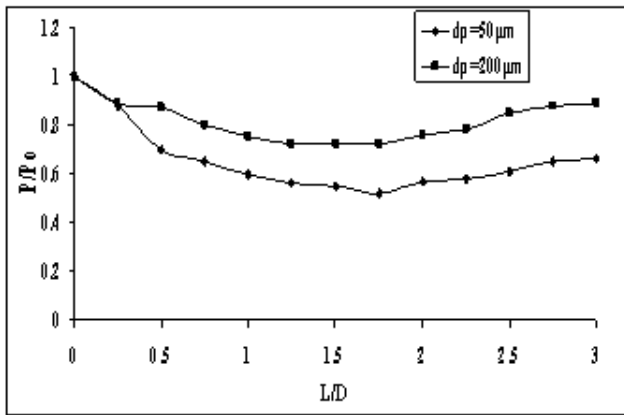
شكل ١ يبين شكل الجهاز المستخدم في البحث



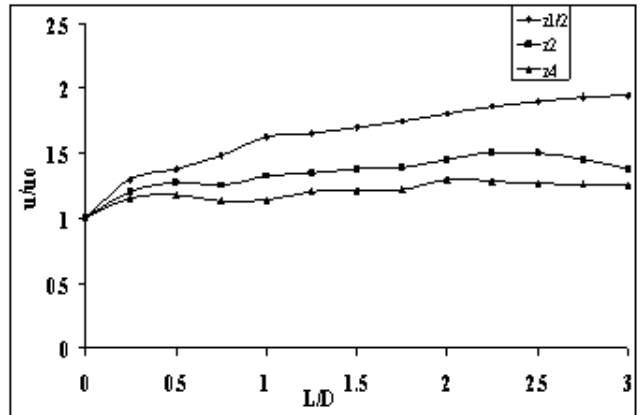
شكل ٣ تغير ضغط الجريان على طول محور الجريان
ولعدة نسب تحميل ولفطر ($m\mu 100$) للجسيم



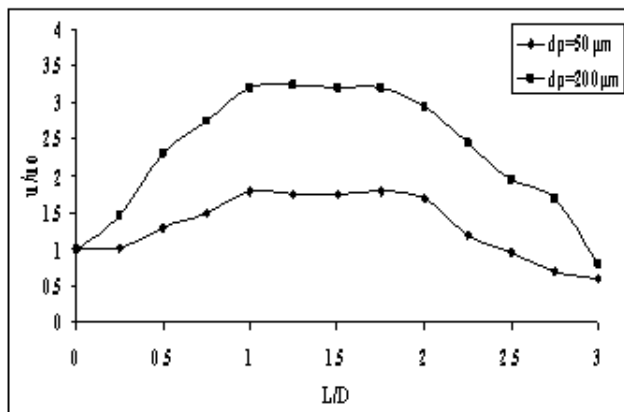
شكل ٤ سرعة الهواء على طول محور الجريان وولعدة
نسب تحميل ولفطر ($m\mu 100$) للجسيم



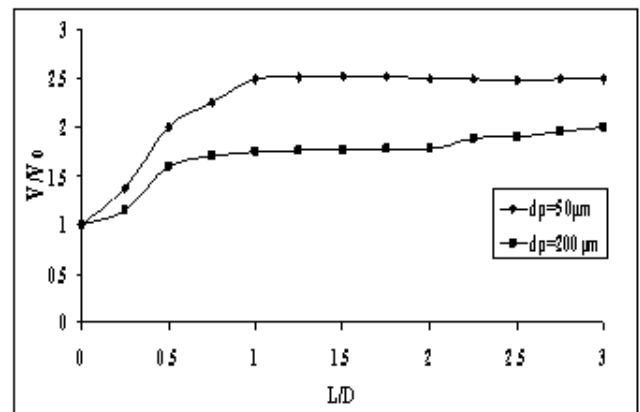
شكل ٥ تغير ضغط الجريان على طول محور الجريان
عند نسبة تحميل (4)



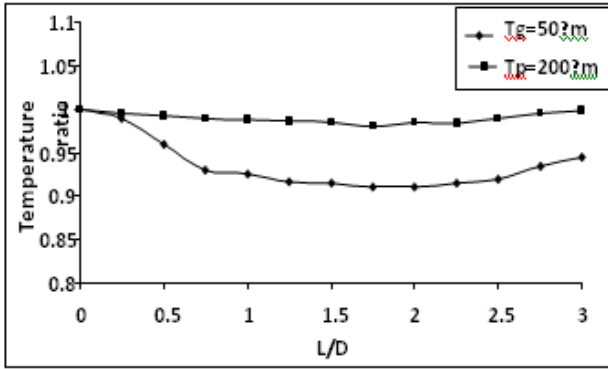
شكل ٦ تغير سرعة الطور الصلب على طول محور
الجريان وولعدة نسب تحميل ولفطر ($m\mu 100$) للجسيم



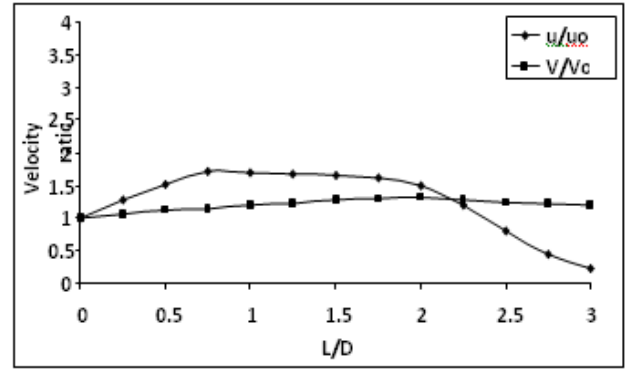
شكل ٧ تأثير حجم الجسيمات الصلبة على سرعة الهواء
عند نسبة تحميل (4)



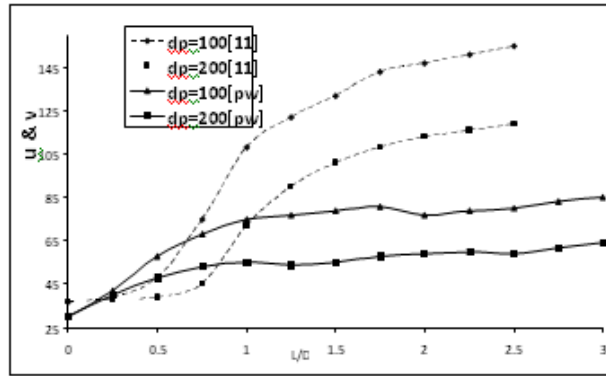
شكل ٨ تأثير حجم الجسيمات الصلبة على سرعة الصلب
عند نسبة تحميل (4)



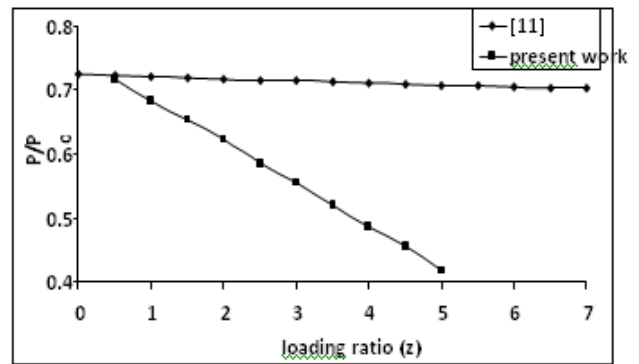
شكل ٩ العلاقة بين درجات حرارة الأقطار عند نسبة تحميل (4) ولقطر (100µm) للجسيم



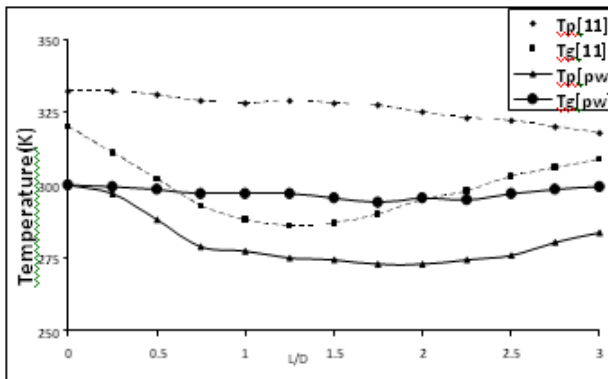
شكل ٨ العلاقة بين سرعتي الأقطار عند نسبة تحميل (4) ولقطر (100µm) للجسيم



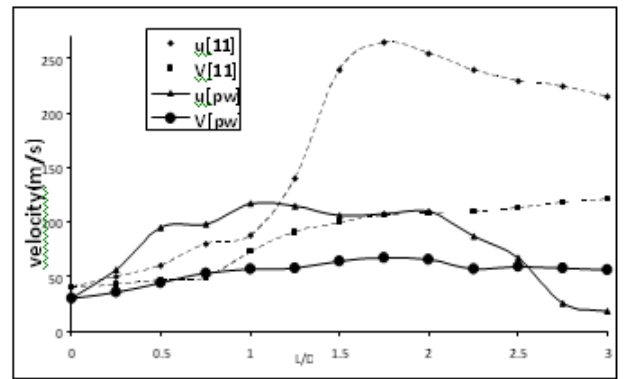
شكل ١١ مقارنة لتأثير حجم الجسيمات الصلبة على سرعة الصلب عند نسبة تحميل (0.5) للعمل الحالي والمصدر [11]



شكل ١٠ مقارنة للعلاقة بين هبوط الضغط ونسبة التحميل للعمل الحالي والمصدر [11]



شكل ١٣ مقارنة لدرجة حرارة الأقطارين عند نسبة تحميل (0.5) للعمل الحالي والمصدر [11]



شكل ١٢ مقارنة لسرعتي الاقطار عند نسبة تحميل (0.5) للعمل الحالي والمصدر [11]